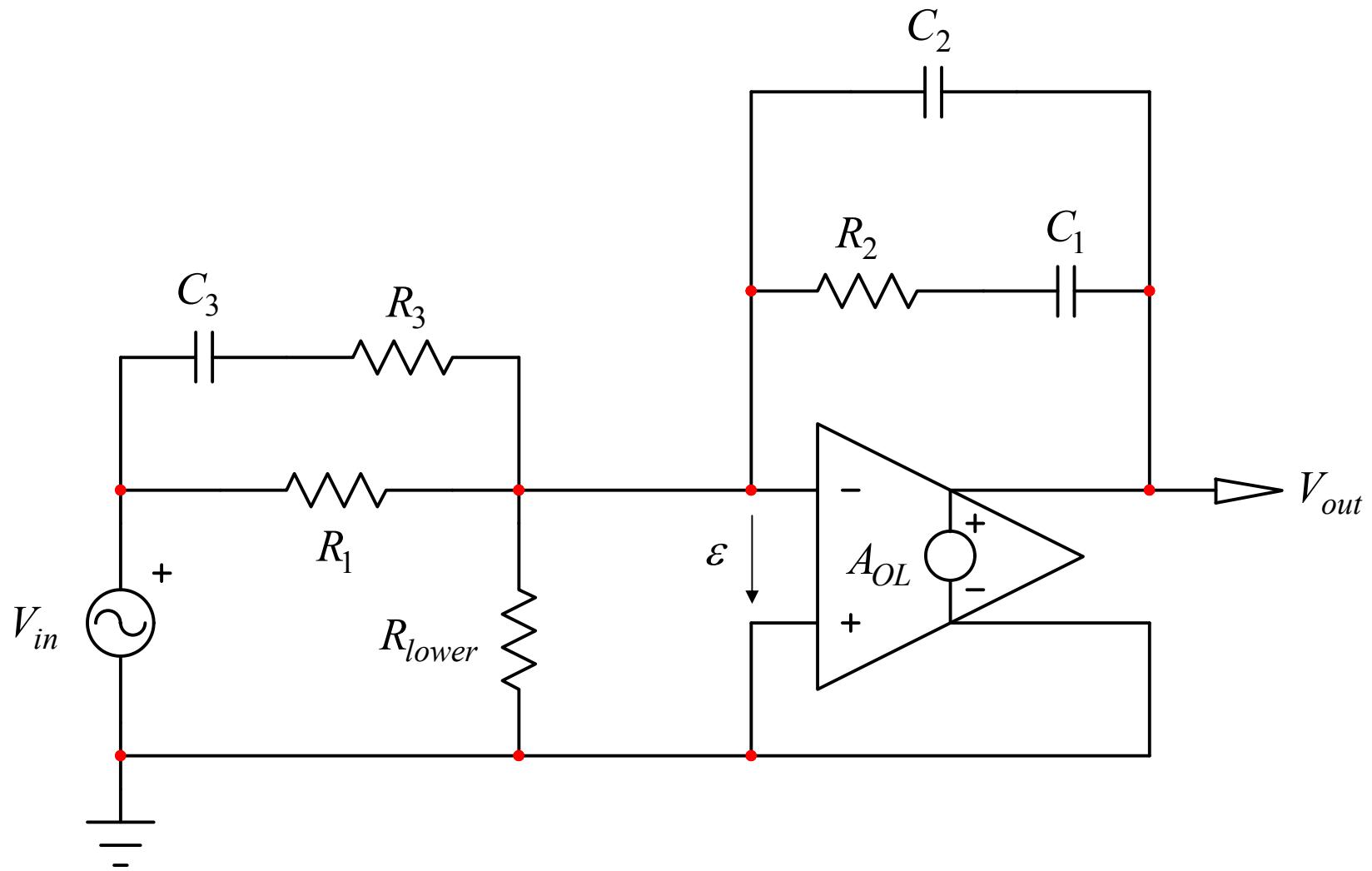
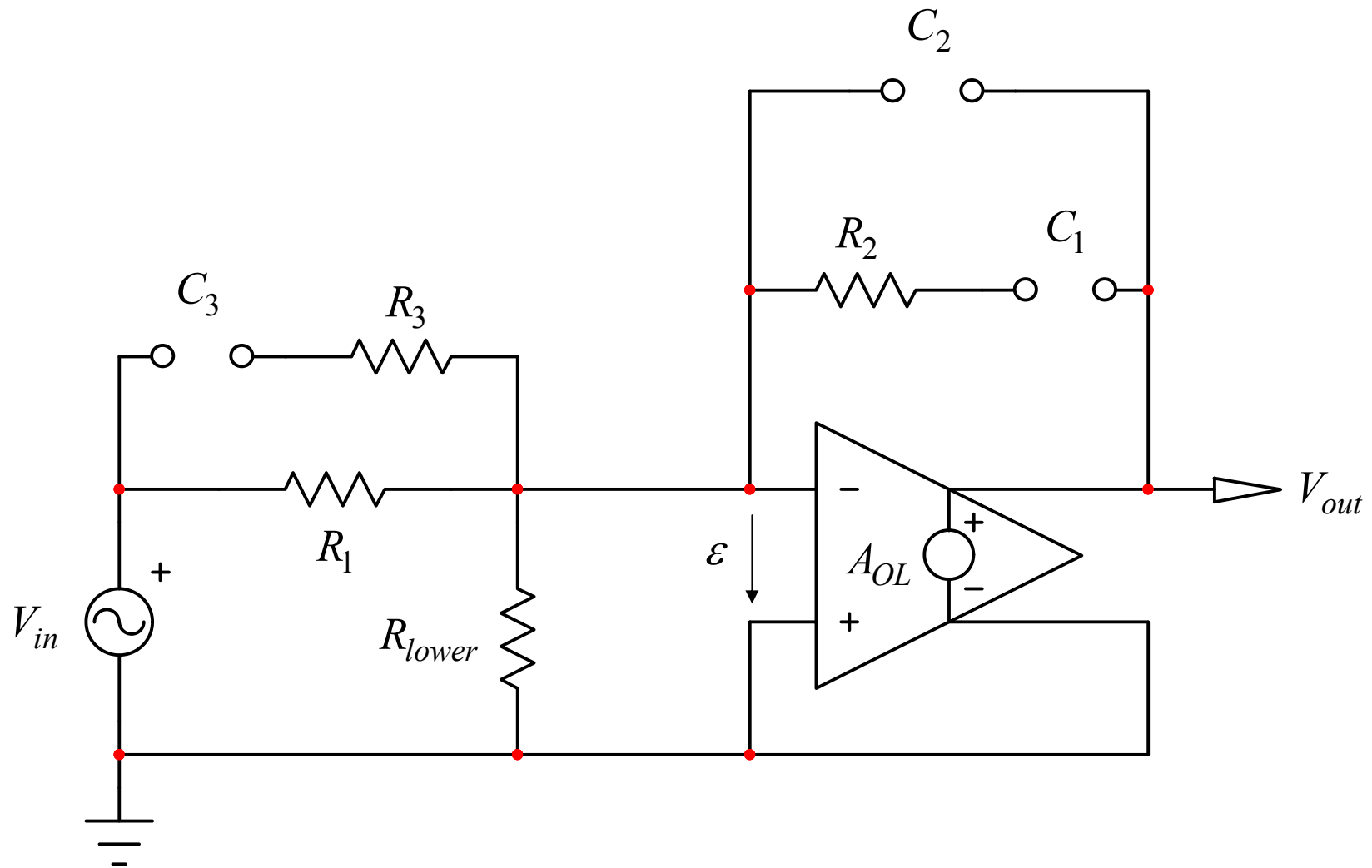


FACTs applied to a type 3 compensator – Christophe Basso – August 2017

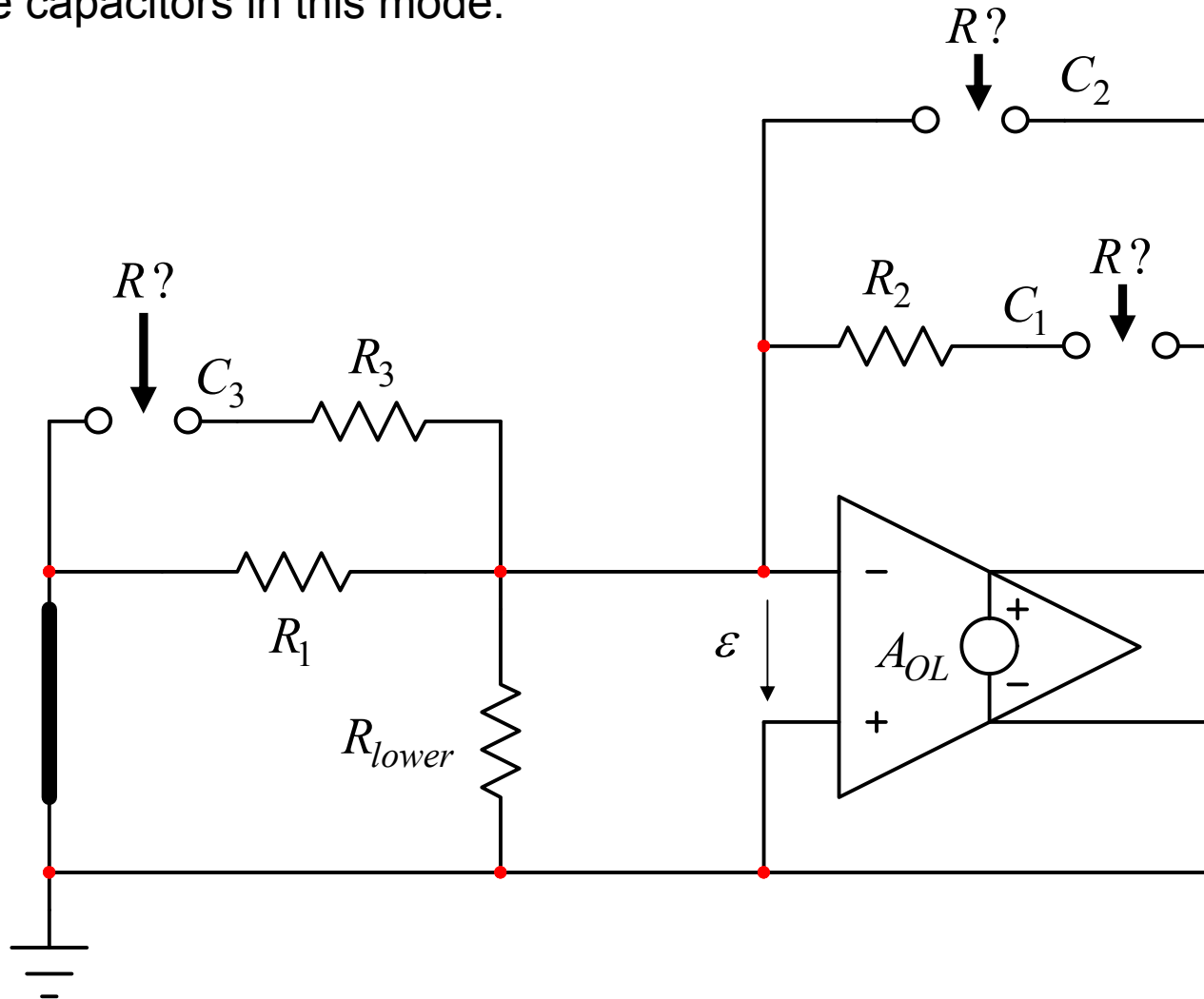


Dc gain, $s = 0$, open all the caps:

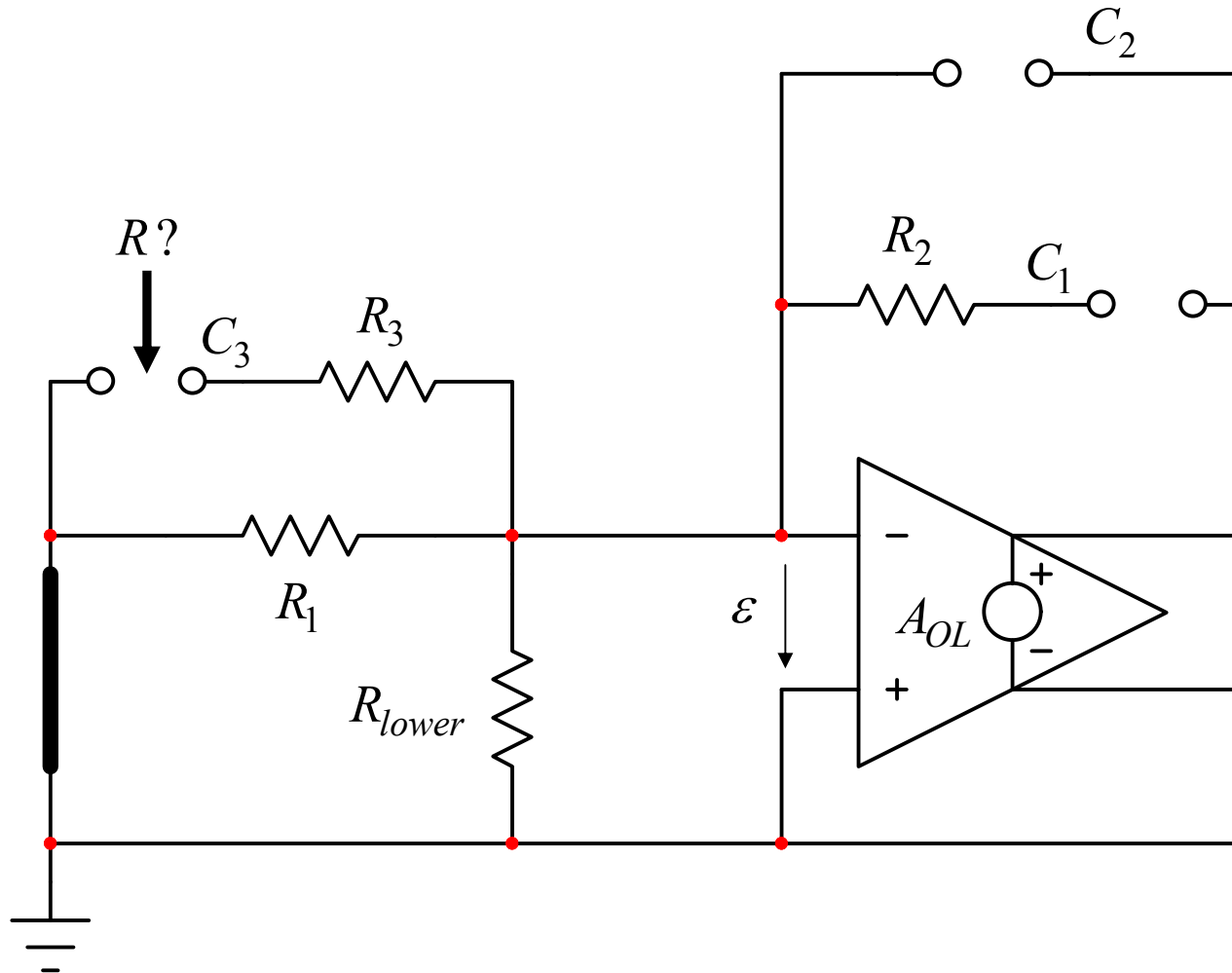


$$H_0 = -\frac{R_{lower}}{R_{lower} + R_1} A_{OL}$$

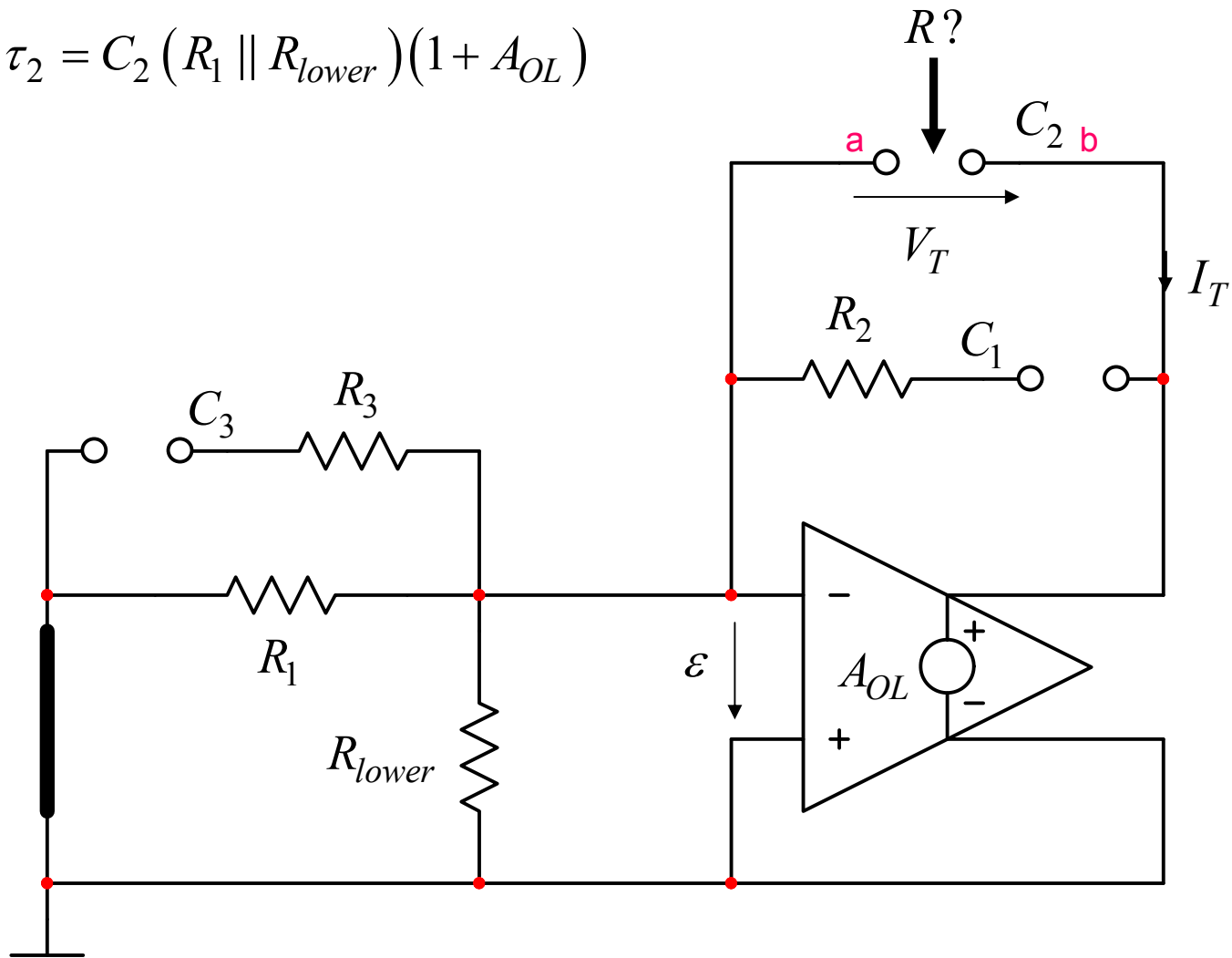
Reduce the excitation (V_{in}) to 0 V and determine the resistances seen from the capacitors in this mode.



$$\tau_3 = C_3 (R_3 + R_1 \parallel R_{lower})$$

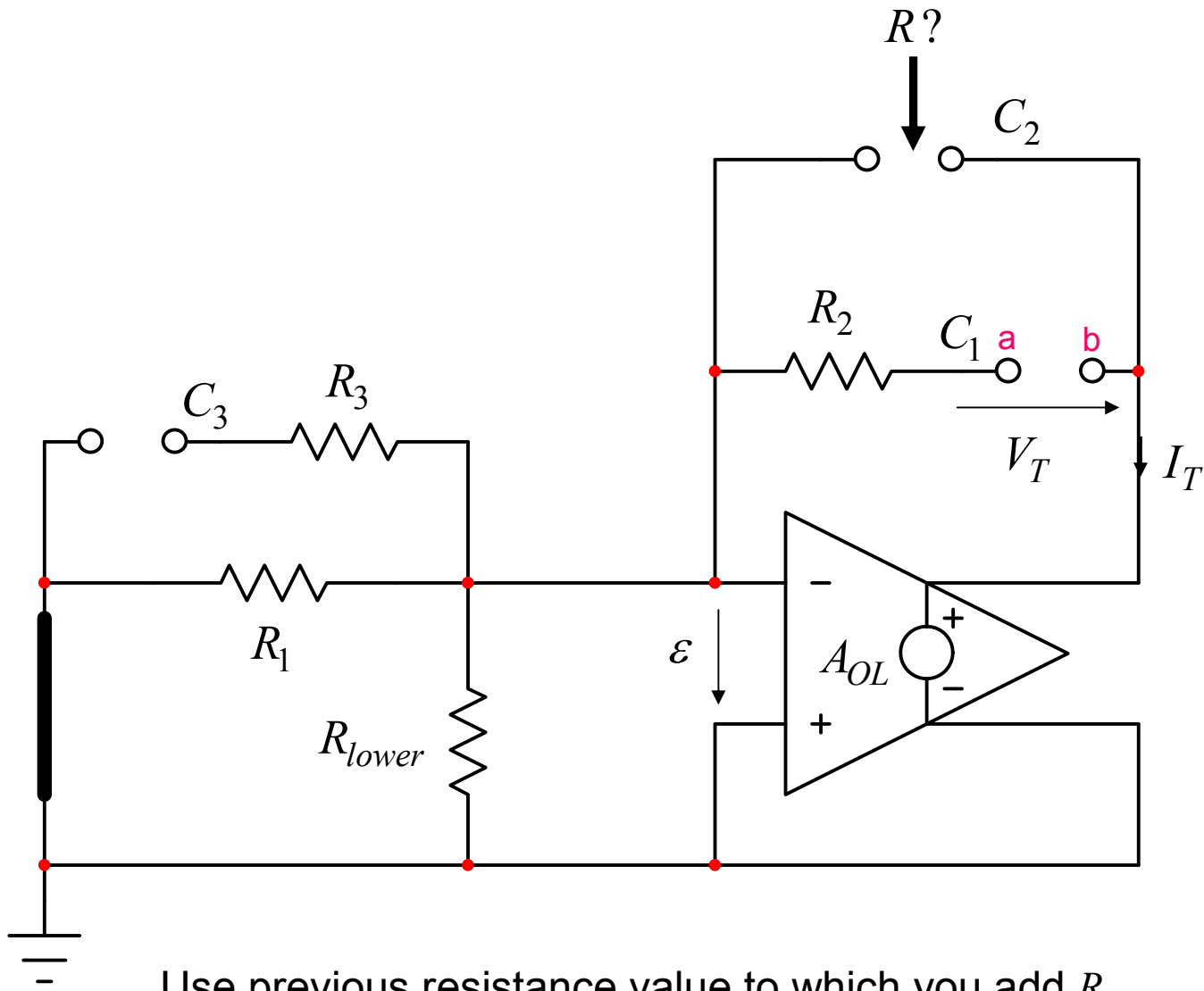


$$\tau_2 = C_2 (R_1 \parallel R_{lower}) (1 + A_{OL})$$



$$V_T = V_{(b)} - V_{(a)} = \epsilon A_{OL} - (-\epsilon) = \epsilon (1 + A_{OL})$$

$$\epsilon = I_T (R_1 \parallel R_{lower})$$



Use previous resistance value to which you add R_2

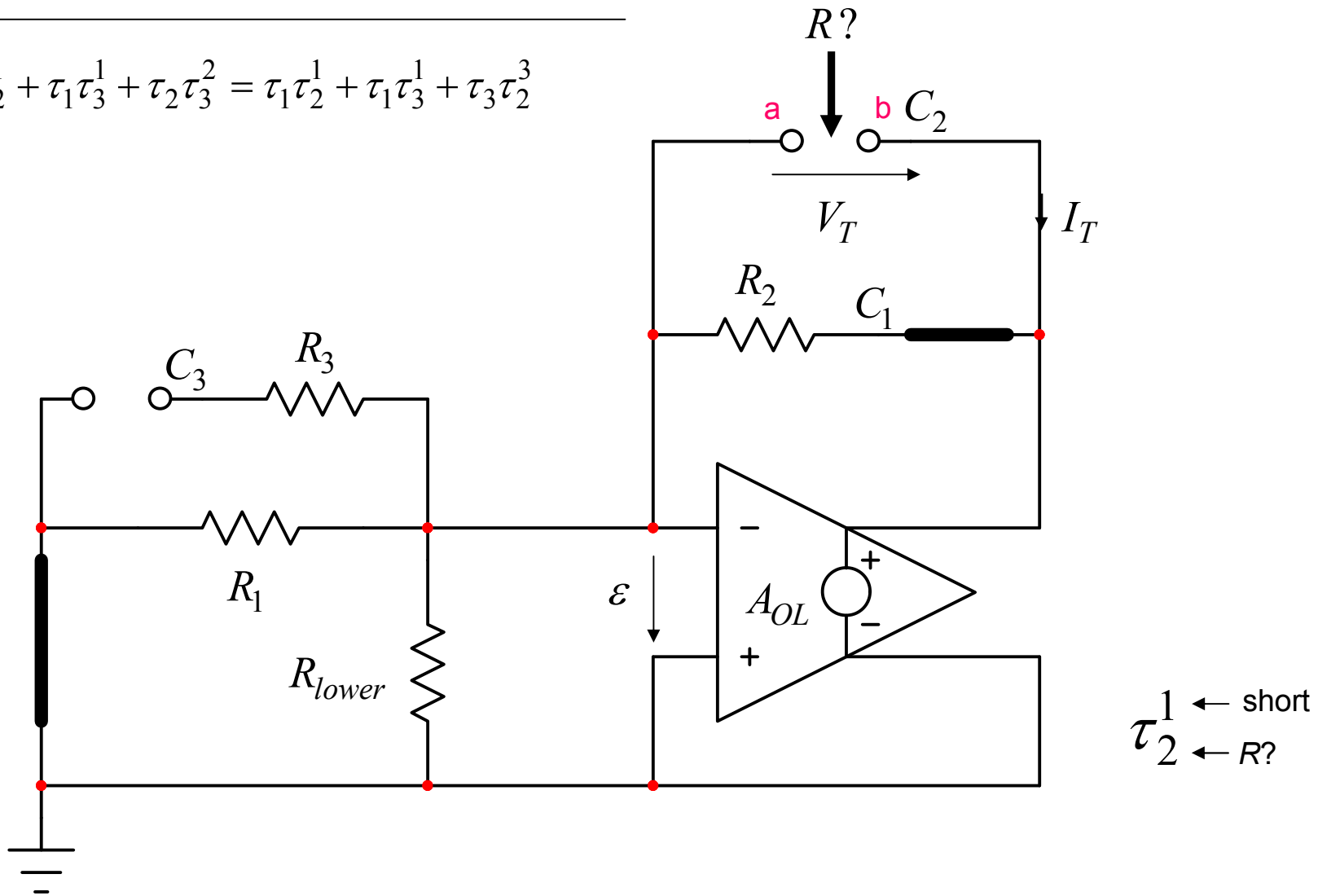
$$\tau_1 = C_1 \left[(R_1 \parallel R_{lower}) (1 + A_{OL}) + R_2 \right]$$

You can form b_1 for the denominator

$$b_1 = \tau_1 + \tau_2 + \tau_3$$

$$= C_1 [(R_1 \parallel R_{lower})(1 + A_{OL}) + R_2] + C_2 (R_1 \parallel R_{lower})(1 + A_{OL}) + C_3 (R_3 + R_1 \parallel R_{lower})$$

$$b_2 = \tau_1 \tau_2^1 + \tau_1 \tau_3^1 + \tau_2 \tau_3^2 = \tau_1 \tau_2^1 + \tau_1 \tau_3^1 + \tau_3 \tau_2^3$$

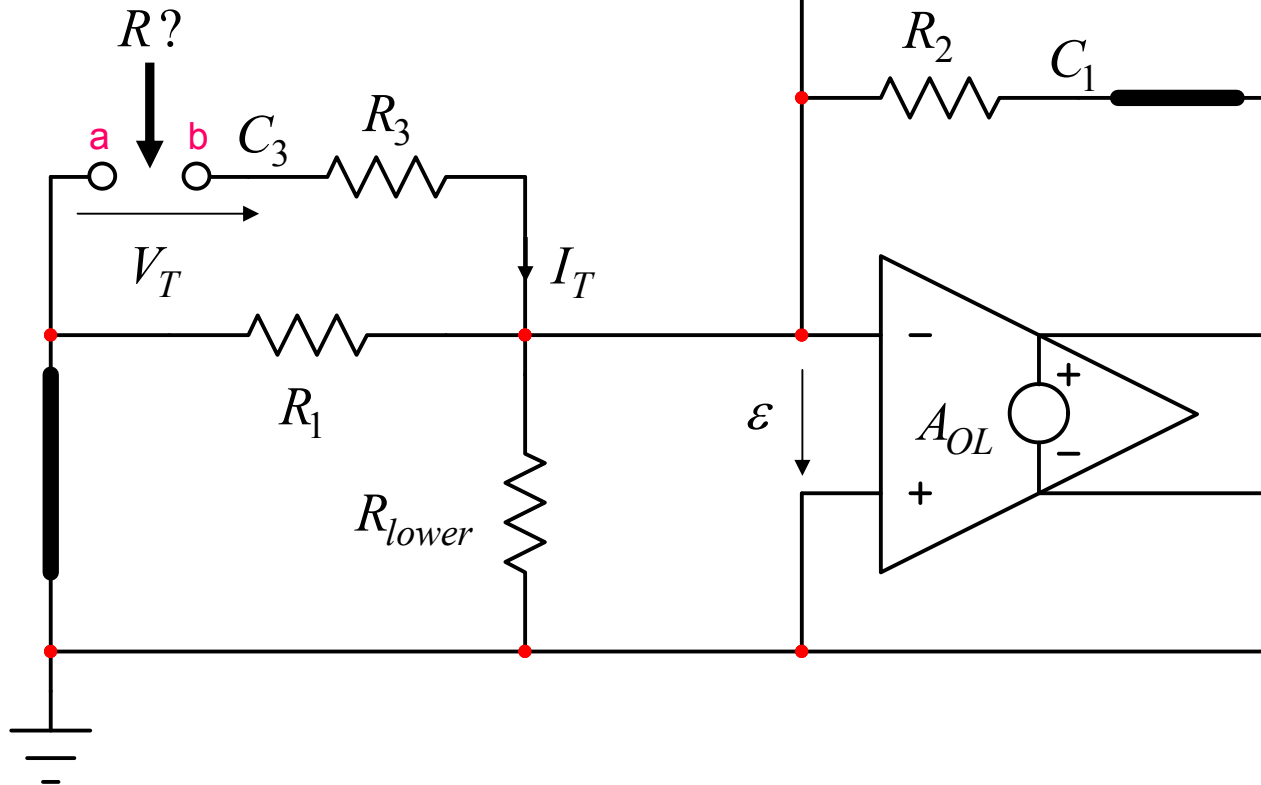


$$V_T = V_{(b)} - V_{(a)} = \varepsilon A_{OL} - (-\varepsilon) = \varepsilon(1 + A_{OL})$$

$$\varepsilon = \left(I_T - \frac{V_T}{R_2} \right) (R_1 \parallel R_{lower})$$

$$\tau_2^1 = C_2 \left[\frac{R_2 (R_1 \parallel R_{lower}) (1 + A_{OL})}{R_2 + (R_1 \parallel R_{lower}) (1 + A_{OL})} \right]$$

$$\tau_3^1 = C_3 \left[\frac{R_2 (R_1 \parallel R_{lower})}{R_2 + (R_1 \parallel R_{lower}) (1 + A_{OL})} + R_3 \right]$$

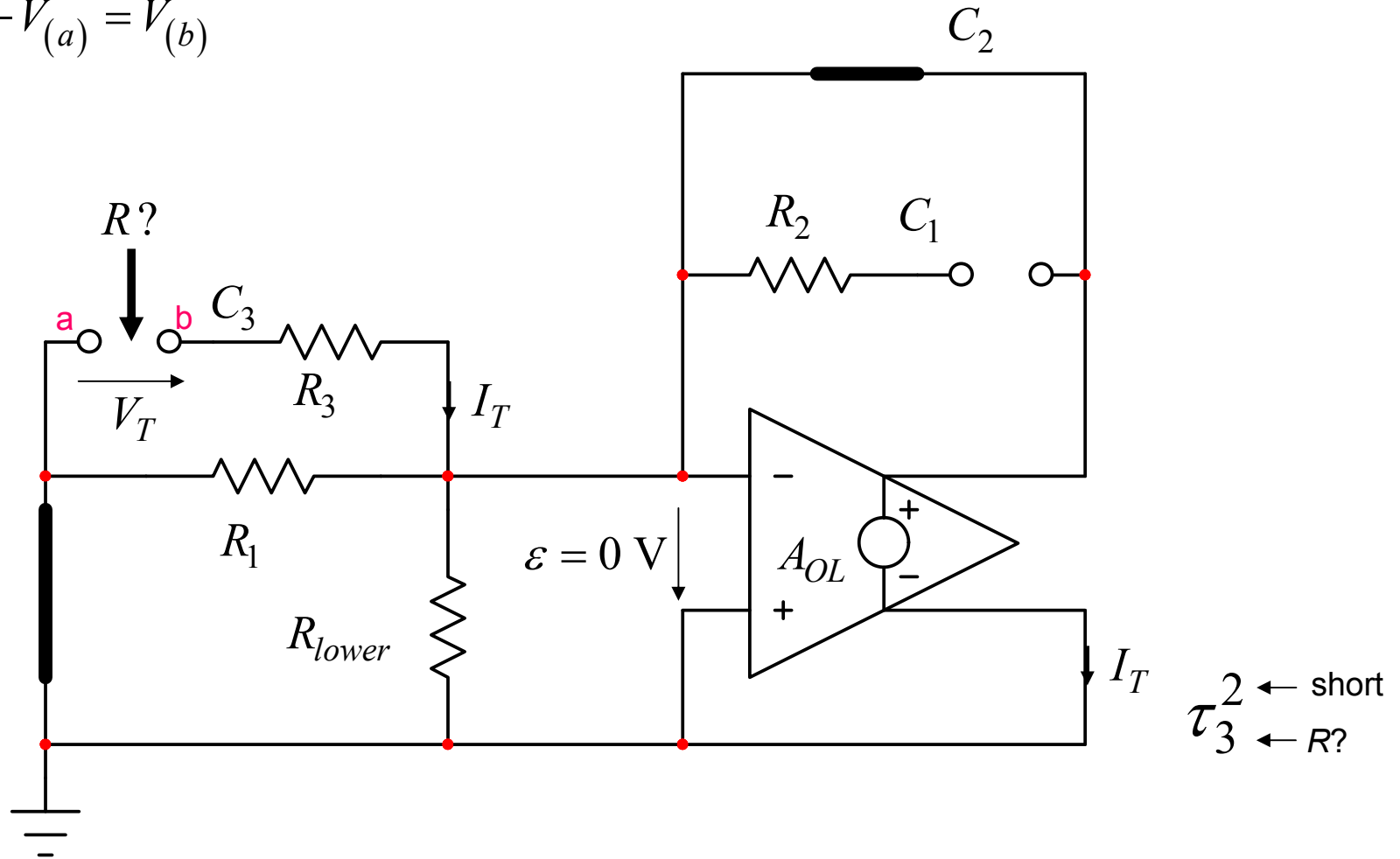


τ_3^1 ← short
 τ_3 ← $R?$

In this configuration, I_T crosses the op amp but not R_1 and R_{lower} .

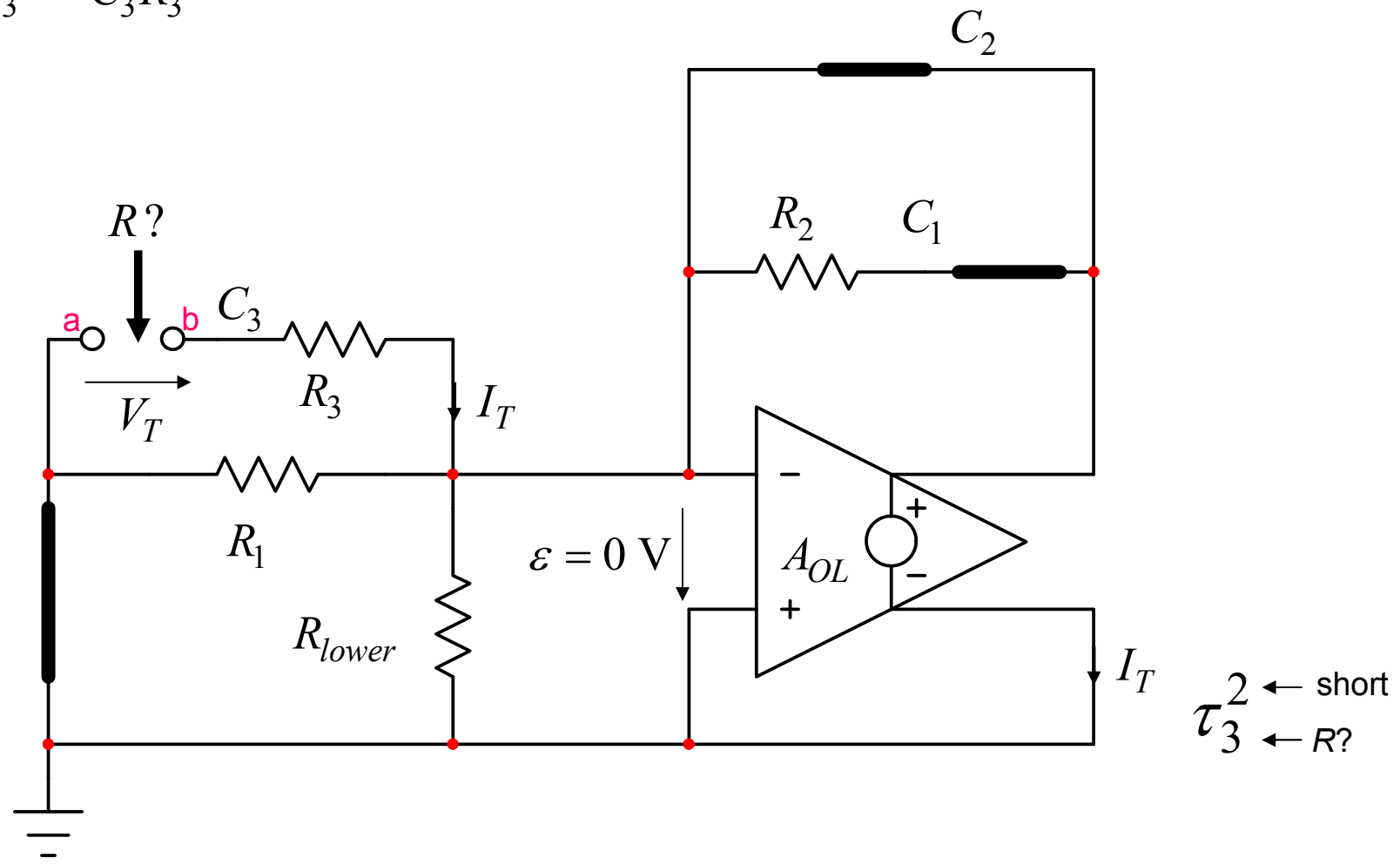
$$V_T = V_{(b)} - V_{(a)} = V_{(b)}$$

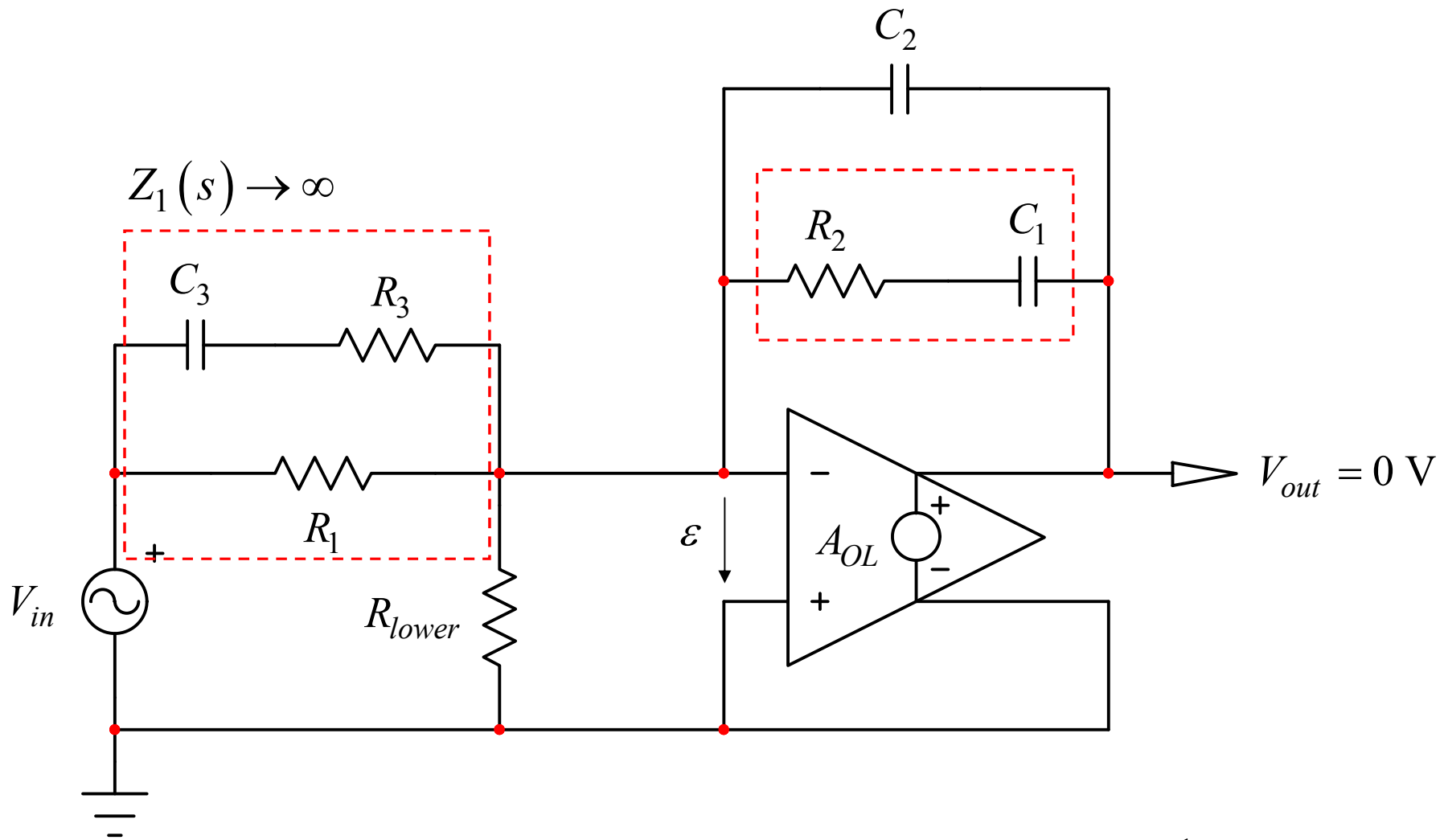
$$\tau_3^2 = C_3 R_3$$



$$b_3 = \tau_1 \tau_2^1 \tau_3^{12}$$

$$\tau_3^{12} = C_3 R_3$$





$$\left. \begin{array}{l} Z_1(s) \rightarrow \infty \\ \frac{N(s)}{D(s)} \rightarrow \infty \\ D(s) = 0 \end{array} \right\} \begin{array}{l} D(s) = 1 + sC_3(R_3 + R_1) \\ \Downarrow \\ \omega_{z_1} = \frac{1}{C_3(R_3 + R_1)} \end{array} \quad \begin{array}{l} Z_2(s) = 0 \rightarrow R_2 + \frac{1}{sC_1} = 0 \\ \Downarrow \\ \omega_{z_2} = \frac{1}{C_1 R_2} \end{array}$$

The final expressions is thus:

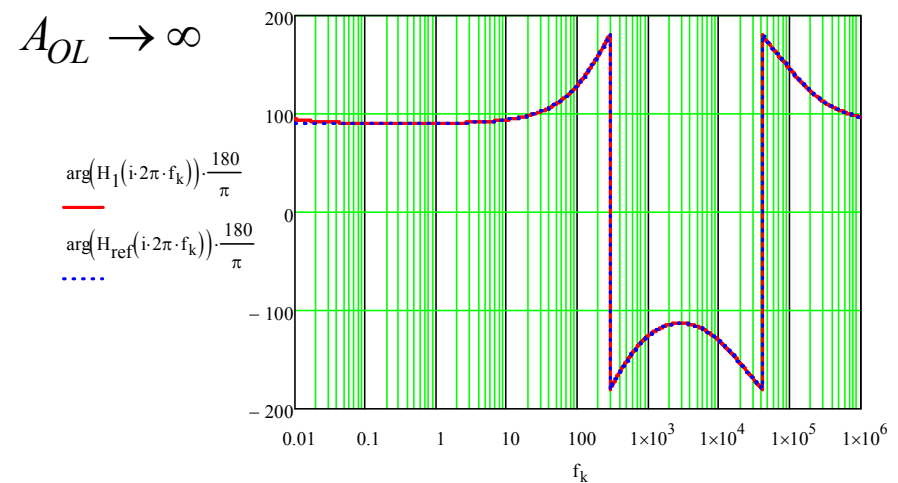
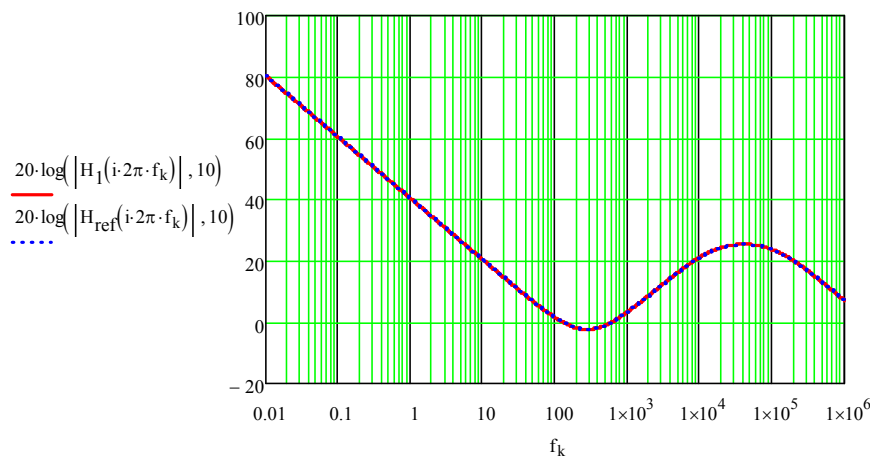
$$H(s) = H_0 \frac{\left(1 + \frac{s}{\omega_{z_1}}\right) \left(1 + \frac{s}{\omega_{z_2}}\right)}{1 + b_1 s + b_2 s^2 + b_3 s^3}$$

$$H_0 = -\frac{R_{lower}}{R_{lower} + R_1} A_{OL}$$

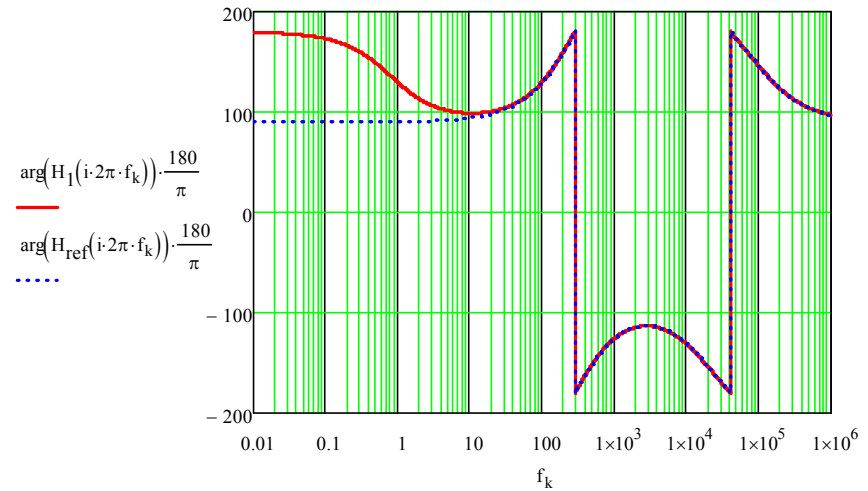
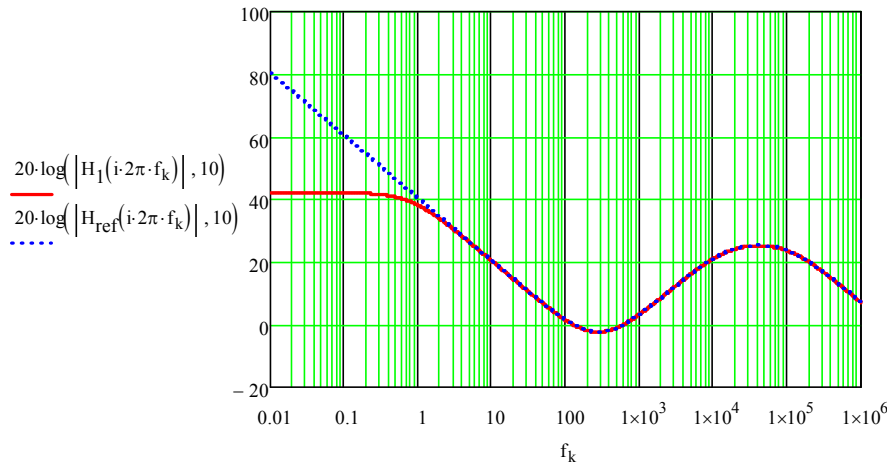
$$\omega_{z_2} = \frac{1}{C_1 R_2} \quad \omega_{z_1} = \frac{1}{C_3 (R_3 + R_1)}$$

The simplified expressions is:

$$H(s) = -\frac{R_2 C_1}{R_1 (C_1 + C_2)} \frac{1 + \frac{1}{s R_2 C_1}}{1 + s R_2 \frac{C_1 C_2}{C_1 + C_2}} \frac{1 + s C_3 (R_1 + R_3)}{1 + s R_3 C_3}$$



$$A_{OL} = 60 \text{ dB}$$



Considering the open-loop gain helps figuring out the role of the resistive network which, in the simplified analysis, does not play a role.

$$H(s) = H_0 \frac{\left(1 + \frac{s}{\omega_{z_1}}\right) \left(1 + \frac{s}{\omega_{z_2}}\right)}{1 + b_1 s + b_2 s^2 + b_3 s^3} \approx - \frac{\frac{R_{lower}}{R_{lower} + R_1} A_{OL}}{b_1 \omega_{z_1}} \frac{\left(1 + \frac{\omega_{z_1}}{s}\right) \left(1 + \frac{s}{\omega_{z_2}}\right)}{\left(1 + \frac{s}{Q\omega_0}\right) \left(1 + \frac{sQ}{\omega_o}\right)} \quad \begin{aligned} Q &= \frac{\sqrt{b_1 b_3}}{b_2} \\ \omega_o &= \sqrt{\frac{b_1}{b_3}} \end{aligned}$$